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# Broadband reflective cholesteric liquid crystalline gels: volume distribution of reflection properties and polymer network in relation with the geometry of the cell photopolymerization

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The ultraviolet (UV) light-absorbing properties of the liquid crystal (LC) constituent during the photo-induced elaboration of a cholesteric LC (CLC) gel may induce the broadening of the reflection bandwidth of the material, a situation that is promoted by asymmetrical irradiation conditions (only one side of the cell is irradiated). The in situ structure of the polymer network, included in the LC, was investigated by transmission electron microscopy and the temperature dependence of the reflection properties examined; it is shown that the network has a structure gradient that is at the origin of the broadening phenomenon. The smallest reflection wavelength is related to the cell side from which the UV light beam came in. A priori, this situation was unexpected since it is shown that this part of the gel is enriched with nematic (infinite-pitch CLC) network-forming material. The result is discussed in relation to the variation of the reflection band characteristics with polymer concentration, which offers the opportunity for indirect access to the volume distribution of the cholesteric periodicities. For applications, broadband reflective cholesteric gels may be of interest for reflective polarizer-free displays or for the light management with smart electrically-switchable reflective windows.

#### 1. Introduction

### 1.1. Property of selective light reflection by a cholesteric liquid crystal

Selective light reflection is exhibited by a cholesteric liquid crystal (CLC) slab with a uniformly oriented Grandjean planar texture (i.e. the helix axis is perpendicular to the observation plane) [1, 2]. At normal incidence, the mean reflection wavelength,  $\lambda_0$ , is related to the pitch p and the mean refraction index  $\bar{n}$  by:

$$\lambda_0 = \overline{\mathbf{n}} \mathbf{p}. \tag{1}$$

 $\bar{n}$  is the average of the ordinary  $(n_o)$  and extraordinary  $(n_e)$  refractive indices of the locally uniaxial structure:  $\bar{n} = (n_o + n_e)/2$ . The spectral width of the reflection band  $\Delta \lambda$  is defined as:

$$\Delta \lambda = p \Delta n, \tag{2}$$

where  $\Delta n = n_e - n_o$  is the birefringence. Within  $\Delta \lambda$ , an incident unpolarized or linearly polarized light beam parallel to the helix axis is split into two opposite

### 1.2. State of the art and objectives of the work

For switchable reflective devices, such as electro-optical glazing structures to be used for the dynamical controlling of the electromagnetic radiations, it is required to go beyond the usual limits of the reflection bandwidth of low molar mass CLCs. The choice of polymer-stabilized CLCs (PSCLCs) elaborated in specific ways has offered the opportunity to increase the reflection bandwidth [3–7]. Polymer-stabilized LCs (PSLCs), also called polymer network SLCs (PNSLCs) or LC gels, are formed by ultraviolet (UV)

circularly polarized components, one of which is transmitted whereas the other is reflected. The sense of rotation of the latter one agrees with the helix screw sense. A wavelength out of  $\Delta\lambda$  is simply transmitted. Compared to conventional pigmented colour filters, a CLC slab has the special feature to combine several optical properties in one layer: it is not only a filter but also a reflector and a polarizer. Since  $\Delta n$  is limited to 0.3–0.4 for colourless organic compounds,  $\Delta\lambda$  is limited to a few tens of nm in the visible spectrum, which means that CLC materials are effective for a certain colour and not for the whole visible spectrum.

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light polymerization of a small amount of bifunctional photoreactive monomers, typically between 3 and 10 wt %, dissolved in a LC phase [8, 9]. The first result of forming a polymer network into the LC solvent is the transfer of the mesophase characteristics to the network morphology and structure during the in situ polymerization. As a consequence, several physical characteristics of the original structure, and consequently the texture, may be kept by the introduction of a large surface to volume ratio as compared to the usual orientation induced by the alignment properties of the cell surfaces. Several research groups have focused their attention on PSLCs due to the peculiar properties they present by comparison with low molar mass LCs. Depending on the pitch range, PSCLCs exhibit specific optical and electro-optical properties which give rise to bistable reflective displays or light shutters with normal [10] and reversed modes [11]. From the more recent literature related to PSCLCs, the following results have been reported: the effect of the UV light wavelength on polymer morphology and electro-optical properties of diffraction gratings [12], the reduction of the hysteresis loop of the reverse-mode PSCL cell by increasing the monomer concentration [13] and the use of PSCLCs for polarization-independent tuneable-focus microlens arrays [14].

With the aim of increasing the reflection bandwidth of CLC structures, an earlier solution consisted of the elaboration of a cholesteric layer with a pitch gradient. A UV light absorbing dye was used to create an intensity gradient of the UV light through the thickness of a solid CLC polymer network [15] or a switchable CLC gel [16]. Therefore, the polymerization at the surface the closest to the UV light source is much faster than at the bottom of the layer. By blending a cholesteric diacrylate with a nematic (infinite-pitch cholesteric) monoacrylate, the former component has a twice as high probability as the nematic component to be incorporated in the polymer network [15]; at low intensity, depletion of the chiral compound near the top of the cell generates a pitch gradient through the thickness of the layer. As a complementary possibility for going beyond the usual limits of the reflection band of low molar mass CLCs, the increase of the reflected light intensity, never greater than 50% of ambient (unpolarized) light, may also be achieved, as recently demonstrated by using PSCLCs with a thermallyinduced inversion of the helicity sense [17, 18]. The present work deals with the broadening of the reflection band of PSCLCs, without the intervention of a foreign UV light-absorbing dye, but as a consequence of the natural absorbing properties of the LC constituent of the gel. We will highlight this phenomenon by using

three different sets of irradiation conditions and, more generally, we will show that it is possible to monitor the bandwidth amplitude by playing with the screening effect due to the LC component of the gel. This result was briefly reported by us in a previous letter intended for a broad readership [19]. The scope of the present paper is to give a detailed description of the bandwidth broadening property by presenting direct reflection profiles from one side of the cell to the other one in relation to the geometry of the cell photopolymerization (asymmetrical vs. symmetrical conditions, presence of a UV light filter) and transmission electron microscopy (TEM) investigations of the in situ polymer morphology. Particular attention is given to an unexpected situation, which consists of having the smallest reflection wavelength near the side of the cell that is enriched in nematic network-forming liquid crystalline material.

### 2. Experimental

#### 2.1. LC blend

A room-temperature CLC mixture (BL094, Merck Ltd) is mixed with bifunctional photoreactive nematic mesogens (RM257, Merck Ltd) at a concentration ratio of 91 wt% of LC to 9 wt% monomers, with a small amount of photoinitiator Irgacure 907 (Ciba-Geigy; 2 wt% compared to the RM257 amount). The blend exhibits a cholesteric phase at room temperature. From texture investigations by polarizing optical microscopy (POM), the isotropic phase appears at 75°C.

### 2.2. Cell characteristics

The blend is introduced at  $83^{\circ}C$  by capillarity in an  $8 \, \mu m \, (\pm 0.7 \, \mu m)$  thick ITO-coated glass cell (from EHC Co. Ltd; type FF). In order to induce a planar alignment, the surfaces were treated with rubbed polyimide (the plates were assembled with their rubbing directions making an angle of  $180^{\circ}C$ ). Then the cell was quenched down from  $70^{\circ}C$  to room temperature. When observed by POM in transmittance mode, the cell exhibits a well-aligned reflective Grandjean planar texture with a moderate presence of oily streaks.

### 2.3. UV light curing conditions

The cell was kept at room temperature and, for polymerization purposes, illuminated for 30 min with UV light centred at 365 nm and with a power equal to 0.1 mW cm<sup>-2</sup> as measured with a radiometer (UVR-365, Prolabo). In the following, the so-called *asymmetrical conditions* correspond to an irradiation from one single side (with a black substrate disposed under the

cell), whereas the *symmetrical conditions* correspond to an irradiation from both sides simultaneously.

### 2.4. Spectrophotometry

The transmission and reflection spectra were recorded by using a HR2000CG-UV-NI spectrophotometer (Ocean Optics) linked to an optical microscope. When measurements in the infrared spectrum are required (see figures 8–9), the UV-3100 spectrophotometer (Shimadzu) was used for all the measurements. The incident light was unpolarized. All the baselines were recorded with air (a mirror) in the beam path for the transmittance (reflectance) measurements. The multiple secondary peaks present in the transmittance spectra correspond to the wave interference due to the internal reflection between the two inner surfaces of the cell substrates; the periodicity is related to the film thickness.

### 2.5. Preparation of materials for TEM and observation conditions

Resin was substituted for the LC component by immersing the gel cell in a bath of LR White Resin diluted in a solvent (methanol); the concentration of resin in the solvent was progressively increased from one bath to the next to optimize the filling of the network porosity. Such a *modus operandi* was expected to preserve the spatial distribution of the network within the cell volume. Finally, the cell was opened at room temperature, the material, which consists of the polymer network embedded in a resin lump, removed

and embedded in a second (epoxy) resin in order to be manipulated. 80 nm thick slices of cross-sections of the sample were cut with an ultramicrotome (Reichert) at room temperature. The cuts are observed by using a Philips CM12 transmission electron microscope.

### 3. Results

### 3.1. Polymerization under asymmetrical conditions

Figure 1 shows the transmittance spectra of the cell before and after the photopolymerization. Before curing,  $\Delta\lambda$ , as measured at the half-width of the peak, is about 80 nm. After curing, the reflection bandwidth is approximately located at the same mean wavelength and has been significantly increased:  $\Delta\lambda$  is now equal to about 220 nm.

Figure 2 shows the reflection spectra of the cell after curing according to the side of the cell that is presented to the incoming light beam probe of the microscope—spectrophotometer. For analysis purposes, it is indeed relevant to make discrimination between the side of the cell that was nearest to the UV light source, the so-called UP side, and the other one, the so-called DOWN side. Each reflection curve exhibits a main peak, the position of which depends on the side of the cell.

The main peak is located at smaller (longer) wavelengths when the UP (DOWN) side is presented to the beam. The wavelength related to the main peak assigned to the UP (DOWN) side is about 540 nm (660 nm); when observed by POM in reflection mode, the colour of the UP side is greenish, whereas it is reddish for the DOWN side. Although the multiple

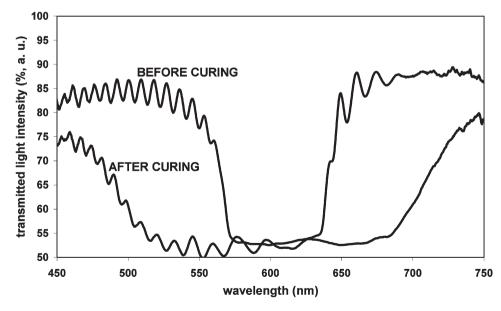


Figure 1. Transmitted light intensity as a function of wavelength before and after curing under asymmetrical conditions.

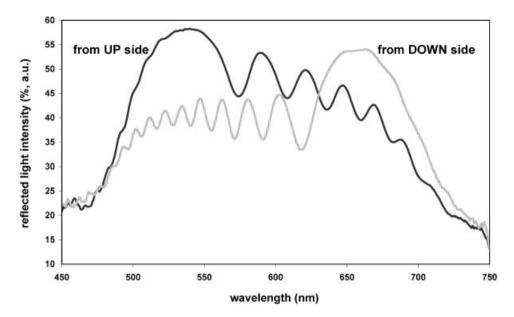


Figure 2. Reflected light intensity related to the UP and DOWN sides as a function of wavelength after curing under asymmetrical conditions.

secondary peaks correspond to the wave interference due to internal reflection between the two inner surfaces of the cell substrates, it appears that colour (Bragg) reflections with a weaker light intensity are also present at wavelengths between the positions of the two main peaks; the intensity is variable and lower in this part of the curves due to a variable and non-optimum number of pitch lengths present in the material [20]. This first set of results suggests that the reflection properties are not monotonously distributed inside the volume of the gel, a

situation from which a broadening of the reflection band may arise, and that the pitch of the hypothetical graded helical structure is longer close to the DOWN side than close to the UP one. The discussion will be focused on this key result.

In order to investigate the relation between the different reflection properties of the two sides, and the possible existence of a concentration gradient, we studied the variation of the reflection spectra with the temperature. Figure 3 shows two sets of results

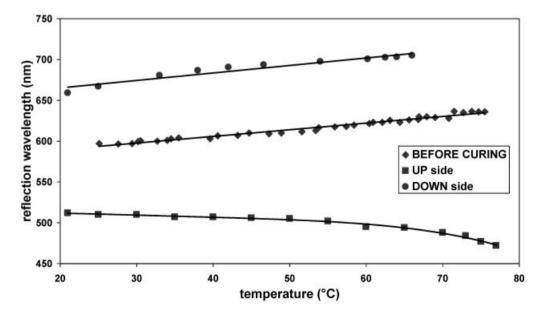


Figure 3. Reflection wavelength related to the UP and DOWN sides as a function of temperature after curing under asymmetrical conditions. The data related to the blend before curing are also reported.

depending on whether the UP or DOWN side of the cell is presented to the incoming light beam. The variation of the reflection wavelength with temperature for the blend before curing is also reported for comparison purposes; in this case, as expected, the reflection property stops at the clearing temperature (above 75°C).

A reflection wavelength assigned to the UP side is not reported above 78°C because the measurements were difficult due to the weak quantity of the related light intensity (and the limit of sensitivity of the spectrophotometer); however, it is important to notice that cholesteric planar textures may be observed by POM (reflection mode) up to 150°C and above, like the consequence of the patterning effect of the polymer network located close to the UP side. PSLCs are indeed commonly described with a two-population model for the low molar mass LC [21]. One population, the free fraction, is weakly influenced by the presence of the polymer, whereas the other one, the bound fraction, undergoes strong surface interactions with it. The bound fraction corresponds to LC molecules in the immediate surrounding of the polymer walls. The polymer morphology ensures memory effects of the orientational order present when its formation occurs (anisotropic growth of polymer units, like, for example, oriented fibres). Therefore, the regions of gel corresponding to the bound fraction will encompass the isotropic transition at a much more elevated temperature in comparison with the behaviour of the free fraction, which has similarities with the behaviour of the LC before its blending with the network-forming material. The report of a reflection wavelength assigned to the DOWN side also stops at 65°C due to the weak quantity of the related light intensity. When the temperature is increased from 20°C up to the clearing point, the wavelength is shifted towards slightly smaller (longer) values for the UP (DOWN) side.

On the one hand, the variations of  $\lambda$  with temperature for the blend and the DOWN side have similar behaviours. The DOWN side is mainly made of a free population of LC molecules (with a fraction of the network as oligomer chains at the very most), a situation that could explain why the variation of  $\lambda$  is mainly dictated by the pitch variation.  $\lambda$  for the DOWN side is slightly but systematically located above the  $\lambda$  attached to the blend. This difference may be understood as the consequence of the composition of the CLC slab close to the DOWN side, which is different from the blend composition since it is deprived of a part of the network-forming material and includes a sparse network (oligomers), as the TEM investigations will show.

On the other hand, the variation corresponding to the UP side is quite different: the reflection wavelength slightly decreases with temperature and close to the clarification temperature of the blend before curing. The UP side is made of a dense and thermally-stable network; the variation of the reflection wavelength would come from the (modest) variation of the mean optical index of the LC component (which decreases up to the isotropic point) included in the gel layer close to the UP side more than the pitch variation because the pitch is here locked in the network.

Thus, the results displayed in figure 3 again suggest that the reflection wavelength is greater close to the DOWN side than close to the UP one and that the polymer concentration is larger close to the UP side.

### 3.2. Polymerization under symmetrical conditions

Since it is known that a broadening of the reflection band may arise from a UV light intensity gradient [15], we investigated the role of a dual-irradiation on the broadening property of the reflection bandwidth. We thus arranged a cell at half-distance between the two UV light sources. A mean power equal to 0.1 mW cm<sup>-2</sup> was measured on each plate of the sandwich-cell. In that case, the phenomenon of broadening of the reflection band is still present but less pronounced in comparison with the case of asymmetrical conditions: the reflection bandwidth is equal to about 160 nm vs. 220 nm in the former case, as reported elsewhere [19]. The reflection spectra for both sides of the cell are displayed in figure 4: the peaks are now very close and located around 620 nm.

The variations of the reflection wavelength with the temperature now have very similar behaviours for both sides (figure 5).

They are located slightly above the curve assigned to the variation of the reflection wavelength before curing; however, the difference is very small and the relative position of curves, as well as their position with regard to the curve assigned to the blend, have no clear signification since the difference between values reported in figure 5 are comprised in measurement errors. Thus, the asymmetry in the position of the reflection band is lost in the case of a UV light irradiation from both sides simultaneously. Under the hypothesis of a concentration gradient in polymer network at the origin of the pitch gradient, we can understand that this gradient will be modest, although still present, when the symmetry for the fabrication of the network in the LC volume is met. This result is evidenced by the proximity of curves in figure 5 when we compare it to the dispersion of measurements in figure 3. The fact that the reflection property of both

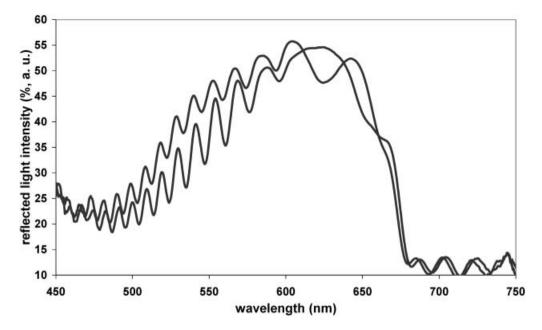


Figure 4. Reflected light intensity related to both sides of the cell as a function of wavelength after curing under symmetrical conditions.

sides is close to the one of the blend before curing shows that the quite homogeneous distribution of the cholesteric periodicity initially present in the CLC layer has been almost frozen in the gel structure and that the in situ formation of the polymer network does not have a significant impact on the reflection property when the cell is irradiated on both sides at the same time.

## 3.3. Polymerization under asymmetrical conditions in the presence of a UV light filter made with the LC component

Since the broadening of the reflection band might be related to a concentration gradient in the polymer morphology from the UP side of the cell to the DOWN one, a progressive absorption of UV light by the LC

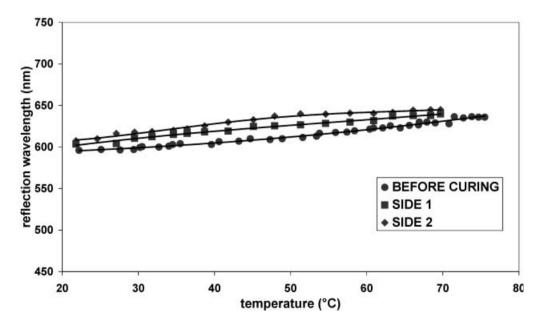
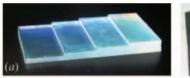


Figure 5. Reflection wavelength related to both sides of the cell as a function of temperature after curing under symmetrical conditions. The data related to the blend before curing are also reported.





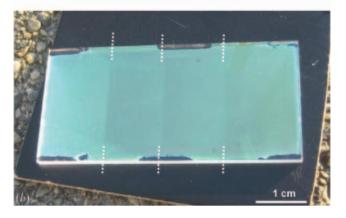


Figure 6. (a) The staircase-like filter: a thickness-gradient cell filled with BL094; lateral (left) and top (right) views are shown. (b) Large-size  $(2.6 \times 5.2 \, \text{cm})$  homemade cell after curing when the staircase-like filter was attached to one side during the UV light irradiation under asymmetrical conditions. The frontier zones corresponding to the interface between two successive steps of the filter (now removed) are visible; they are partially stressed with white dotted lines in the photograph.

constituent may be suspected. We then decided to cure the experimental cell when a cell filled with the pure LC (BL094) is attached to it (i.e. the UP side). Such geometry offers the opportunity to investigate the screening role of the LC constituent on the UV light beam and, in return, on the polymerization characteristics and the optical properties. The variation of the bandwidth with the thickness of the LC slab used like a filter was investigated.  $\Delta\lambda$  linearly decreases when the thickness increases; whereas  $\Delta\lambda$  is about 220 nm when the curing occurs without filter, it decreases from 180 to 100 nm when the filter thickness changes from 4 to 44 µm, as reported previously [19]. This result shows that the part of the incoming UV light beam has been filtered by the LC cell and that the screening effect, from which a gradient network may arise, is strongly attenuated. Figure 6 provides macroscopic evidence of this result by displaying a photograph of a large-size  $(2.6 \times 5.2 \,\mathrm{cm})$  homemade cell filled with the blend when the curing occurs with a staircase-like filter cell (figure 6 a) attached to it.

In figure 6a, one, two, three and four LC slabs of BL094 are successively present from the left side of the filter to the right one; each LC slab is about 10 µm thick. A variation of the green reflection colour can be observed in a direction parallel to the long side of the experimental cell (figure 6b). The frontier zones, which

are partially stressed with white dotted lines in figure 6 b for presentation purposes, mark the limit between two successive steps of the staircase-like filter (which has been removed here).

### 3.4. Investigations of the polymer network distribution and morphology

TEM investigations offer a valuable opportunity to obtain information about the structure of the polymer network in relation with the UV light curing conditions. Figure 7 shows the micrographs of cross-sections of the polymer when the gel was fabricated under asymmetrical (figure 7 a) or symmetrical (figure 7 b) conditions. Figure 7 a clearly shows that the network is more concentrated close to the UP side of the cell. A scattered network is displayed in the neighbourhood of the opposite side of the cell: this situation suggests that a discontinuous network, with oligomeric structures, is formed close to the DOWN side. Therefore, a concentration-gradient network is formed in the LC volume when the curing conditions are asymmetrical. Under the hypothesis of a concentration gradient in the polymer network as the origin of the pitch gradient, we can understand that this gradient is more modest when the symmetry conditions for the fabrication of the network in the LC volume are met and why a less broadened bandwidth is obtained in return. Indeed, figure 7 b shows that the network morphology and distribution are qualitatively the same from one side to the other one by being symmetrical around the 'symmetry' axis as reported in the image. Interestingly, it has to be noticed that a discrete concentration gradient is still present from each one of the surfaces. In both cases, a local orientation of chains of polymer beads may be distinguished in a direction parallel to the surfaces, but this kind of organization does not appear to be dense and regular, whereas we observed a periodic stratification reminiscent of the cholesteric fingerprint texture in 25 µm thick samples. These images are the first TEM visualization of the in situ distribution of a concentration-gradient polymer network in the volume of a PSLC.

### 4. Discussion

We have shown that a broadening of the reflection bandwidth in a PSCLC may arise from the formation a concentration-gradient polymer network, which is promoted by asymmetrical curing conditions as a consequence of UV light screening by the LC component. The polymer concentration is greater close to the UP side than close to the DOWN one. Consequently, since the network-forming material is a nematic LC

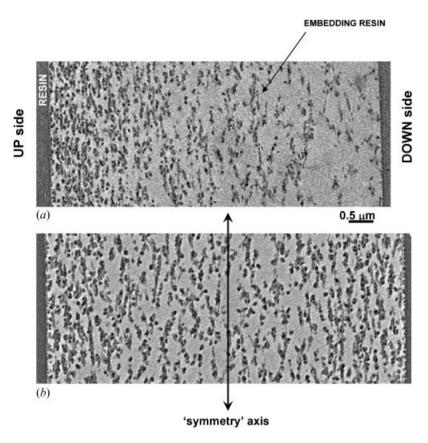


Figure 7. Transmission electron microscopy images of the polymer network embedded in a resin matrix, after dissolution of the LC, when the cell was irradiated: (a) from one side – the UP side (asymmetrical conditions); (b) simultaneously from both sides (symmetrical conditions).

(infinite-pitch CLC), the pitch would have been expected to be longer close to the UP side and decreasing in a direction perpendicular to the UP substrate, as reported in the literature [6, 22]. But this scenario is in apparent contradiction with the present reflection measurements, which correspond to the smallest reflection wavelength close to the UP side. We are going to discuss this point. Since the broadening of the reflection bandwidth comes from a network concentration that decreases from the UP to the DOWN side, it is relevant to investigate the variation of characteristics of the reflection band,  $\lambda$  and  $\Delta\lambda$ , when the concentration in network-forming material varies.

Firstly, from figure 8 (from transmittance measurements and under asymmetrical conditions) we learn that  $\lambda$  varies over a broad range from 540 up to 1800 nm, i.e. from the middle of the visible spectrum up to the infrared spectrum, when the RM257 concentration increases from zero (pure BL094) to 70 wt%. This means that cholesteric periodicities with related  $\lambda$  well above the focused range of the present investigation (450–750 nm) may be present in the helical structure of the gel layer. Besides, the position of the reflection band

is here remarkably not affected by the polymerization process.

Secondly, the investigation of the variation of  $\Delta\lambda$  with the RM257 concentration (from transmittance measurements and under asymmetrical conditions), in figure 9, gives indirect information of paramount importance on the distribution of periodicities to be brought together with the polymer distribution, as seen in figure 7(a), and suggests the following:

- (i) The variation of  $\Delta\lambda$  with the RM257 concentration exhibits a non-monotonous behaviour. Especially,  $\Delta\lambda$  increases in a quite broad range, from about 100 to 420 nm, when the RM257 concentration changes in a restricted range, from 4 to 17 wt %.
- (ii) Then, after an inflexion point, at about  $20 \, \text{wt} \, \%$ ,  $\Delta \lambda$  decreases more slowly from 430 to 165 nm when the concentration varies from 23 to 60 wt %.
- (iii) Finally,  $\Delta\lambda$  increases with a behaviour similar to the blend one from 60 wt%; the network concentration is so large that a distinction between the free and bound fractions of LC molecules has no more sense and the material is constituted of LC

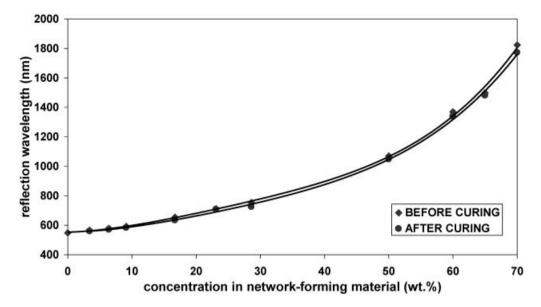


Figure 8. Reflection wavelength as a function of the concentration in network-forming material (RM257) before and after curing under asymmetrical conditions.

molecules closely associated to the network (bound fraction). The fact that  $\Delta\lambda$  is smaller after curing comes from a pitch decrease related to the volume shrinkage resulting from the reaction of polymerization and crosslinking; this behaviour was also observed in the formation of CLC polymer networks [23] and CLC gels [18].

Because the DOWN side is less enriched in polymer network, we could liken the composition of this region of the cell to the first part of the curve in figure 9, i.e. a region for which a large set of reflection bandwidths is present and in relation to the smallest concentrations in network-forming material. It could explain why the reflection wavelength exhibited by the gel when the DOWN side is presented to the incoming beam is greater. As an example, and in comparison with what would happen in the neighbourhood of the UP side, it would be required to have a concentration profile sampled on 35 wt % (from 25 to 60 wt %), vs. 10 wt % for the left part of curve in figure 9 (from 7 to 17 wt %), to reach a performance in common, i.e. to have  $\Delta\lambda$  which varies from 165 to 420 nm. In return, from figure 8 it

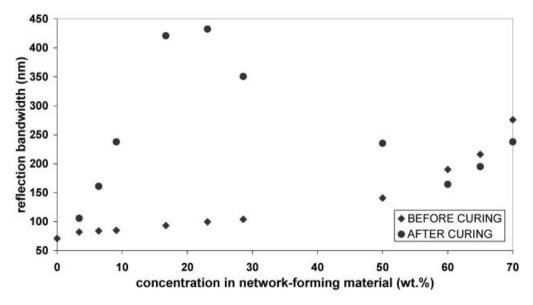


Figure 9. Reflection bandwidth as a function of the concentration in network-forming material (RM257) before and after curing under asymmetrical conditions.

can be seen that the reflection wavelength is out of the investigated range (>750 nm) as soon as the concentration is more than about 28 wt%; an 8 µm thick cell cannot offer the opportunity to let the helix develop a sufficient number of turns at each range of pitch magnitude to give a reflection signal with a significant light flux [20]. Nevertheless, one or several periodicities sampled on all the range exhibited in figure 9 are possible, with a distribution that includes the larger periodicities for the weakest concentrations in networkforming material, close to the DOWN side of the gel layer.

### 5. Conclusion

The natural UV light-absorbing properties of the LC constituent during the photo-induced elaboration of a CLC gel may induce the broadening of the reflection bandwidth, a situation that is promoted when only one side of the cell is irradiated. The polymer component was included in a resin by preserving its spatial distribution and the existence of a structure gradient, which is at the origin of the broadening phenomenon, was shown by TEM investigations of cross-sections. Direct reflection measurements have given clues to the distribution of the reflection wavelength inside the volume of the gel layer. The smallest reflection wavelength was shown to be related to the neighbourhood of the side close to the UV light source when asymmetrical irradiation conditions are met. A priori, this situation was unexpected since the TEM investigations and the temperature-dependence of reflection properties have revealed that this side is enriched in nematic (infinite-pitch CLC) network-forming material. This result was discussed in relation with the variation of the characteristics of the reflection band with the polymer concentration, which offers the opportunity to have an indirect access to the volume distribution of the cholesteric periodicities.

In view of practical applications, broadband reflective cholesteric gels may be of interest for reflective polarizer-free displays or for the light management for smart windows; the latter are electro-optical glazing structures, made with a cholesteric film sandwiched between glass panels, having reflective and transparent modes which are electrically-switchable.

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